

Lecture 4 on 9.4.2024: Subsonic ^{or supersonic} compressible flow around a thin airfoil at zero lift

A) Governing equations for compressible potential flow:

• Momentum + Continuity:

$$(C^2 - \phi_x^2) \phi_{xx} + (C^2 - \phi_y^2) \phi_{yy} - 2\phi_x \phi_y \phi_{xy} = 0 \quad (1)$$

• Total energy + thermodynamic relationships:

$$C^2 = C_\infty^2 + \frac{\kappa - 1}{2} (U^2 - \phi_x^2 - \phi_y^2) \quad (2)$$

where

$$\nabla \phi = \underline{u},$$

$$\kappa = \frac{c_p}{c_v}$$

Non-linear eq. $\rightarrow$ can be solved numerically (FEM / FD)

B) Dimensionless form: - for simplification & generality

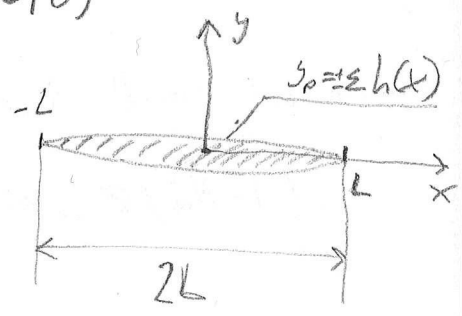
$\rightarrow$ define dimensionless variables:

$$\tilde{x} = \frac{x}{L}, \quad \tilde{u} = \frac{u}{U}, \quad \tilde{C} = \frac{C}{C_\infty}; \quad \tilde{\phi} = \frac{\phi}{UL} \quad \text{s.t.} \quad \tilde{u} = \nabla \tilde{\phi} \quad (3)$$

$\rightarrow$ dimensionless eq.:

$$\left[M_\infty^{-2} + \frac{\kappa - 1}{2} (1 - \tilde{\phi}_y^2) - \frac{\kappa + 1}{2} \tilde{\phi}_x^2 \right] \tilde{\phi}_{\tilde{x}\tilde{x}} + \left[M_\infty^{-2} + \frac{\kappa - 1}{2} (1 - \tilde{\phi}_x^2) - \frac{\kappa + 1}{2} \tilde{\phi}_y^2 \right] \tilde{\phi}_{\tilde{y}\tilde{y}} - 2\tilde{\phi}_x \tilde{\phi}_y \tilde{\phi}_{\tilde{x}\tilde{y}} = 0 \quad (4)$$

c) Recall the problem formulation: $\underline{u}_\infty = (U, 0)$



• Boundary conditions:

— $\underline{\tilde{u}} \rightarrow (1, 0)$ as $\tilde{x}^2 + \tilde{y}^2 \rightarrow \infty$
 $\Rightarrow \tilde{\phi} = \tilde{\chi}$ as $\tilde{x}^2 + \tilde{y}^2 \rightarrow \infty$ (5)
 "far-field"

— symmetry: $\tilde{\phi}_{\tilde{y}} = 0$ at $\tilde{y} = 0$ for $\tilde{x} \leq -1$ or $\tilde{x} \geq 1$ (6)

— kinematic condition:

$\frac{\tilde{\phi}}{\tilde{u}} = \frac{\tilde{\phi}_{\tilde{y}}}{\tilde{\phi}_{\tilde{x}}} = \pm \varepsilon \frac{\partial h}{\partial \tilde{x}}$

 at $\tilde{y} = \pm \varepsilon h$ for $-1 < \tilde{x} < 1$ (7)

d) Asymptotic expansion: (linearization)

— Assumption: the thin profile causes a small perturbation of the incoming flow, proportional to $\varepsilon \ll 1$

$\tilde{\phi} = \phi_0 + \varepsilon \phi_1$ (← drop the $\sim$ for brevity) (8)

$\phi_0 = \tilde{\chi}$ ← undisturbed flow $\underline{u}_0 = (1, 0)$

$\tilde{\phi} = \tilde{\chi} + \varepsilon \phi_1$

$\tilde{\phi}_{\tilde{x}} = 1 + \varepsilon \phi_{1,\tilde{x}}$

$(\tilde{\phi}_{\tilde{x}})^2 = 1 + 2\varepsilon \phi_{1,\tilde{x}} + \varepsilon^2 \phi_{1,\tilde{x}}^2$

$\tilde{\phi}_{\tilde{x}\tilde{x}} = \varepsilon \phi_{1,\tilde{x}\tilde{x}}$

$\tilde{\phi}_{\tilde{y}} = \varepsilon \phi_{1,\tilde{y}}$

$(\tilde{\phi}_{\tilde{y}})^2 = \varepsilon^2 \phi_{1,\tilde{y}}^2 \approx 0$

$\tilde{\phi}_{\tilde{y}\tilde{y}} = \varepsilon \phi_{1,\tilde{y}\tilde{y}}$

$\tilde{\phi}_{\tilde{x}\tilde{y}} = \varepsilon \phi_{1,\tilde{x}\tilde{y}}$

$\Rightarrow$ substitute into governing eq. & BCs
 (drop all $\sim$ for brevity)

D1) Governing equations:

$$\left[M_\infty^2 + \frac{\kappa-1}{2} - \frac{\kappa+1}{2} (1 + 2\varepsilon \vartheta_{1,x}) \right] \varepsilon \vartheta_{1,xx} + \left[M_\infty^2 + \frac{\kappa-1}{2} \cdot (-2\varepsilon \vartheta_{1,x}) \right] \varepsilon \vartheta_{1,yy} - 2(1 + \varepsilon \vartheta_{1,x}) \varepsilon^2 \vartheta_{1,y} \vartheta_{1,xy} = 0$$

- cancel small terms $\varepsilon^2 \ll 1$, $\varepsilon^3 \ll 1$ and divide by ε :

$$\boxed{(1 - M_\infty^2) \vartheta_{1,xx} + \vartheta_{1,yy} = 0} \leftarrow \text{linearized eq. for small disturbances}$$

$\varepsilon \ll 1$ (9)

$\Rightarrow$ can be solved analytically $\rightarrow M_\infty > 1 \Rightarrow$ wave equation

$\rightarrow M_\infty < 1 \Rightarrow$ Prandtl-Glauert transformation
 $\rightarrow$ Laplace eq.

$\Rightarrow$ singularity method

D2) Boundary conditions:

- far-field: $\tilde{\vartheta} = \tilde{X}$ as $\tilde{x}^2 + \tilde{y}^2 \rightarrow \infty$

$$\tilde{X} + \varepsilon \vartheta_1 = \tilde{X}$$

$$\underline{\vartheta_1 = 0}$$

(10)

- symmetry: $\tilde{\vartheta}_y = 0$ at $\tilde{y} = 0$ for $\tilde{x} \leq -1 \vee \tilde{x} \geq 1$

$$\underline{\varepsilon \vartheta_{1,y} = 0}$$

(11)

D2.1- asymptotic expansion of kinematic BC:

$$\phi_y(x, \pm \varepsilon h) = \pm \varepsilon h' \phi_x(x, \pm \varepsilon h) \quad (7)$$

- Taylor expansion: $\phi_y(x, \pm \varepsilon h) = \phi_y(x, 0^\pm) \pm \varepsilon h \phi_{yy}(x, 0^\pm) + \mathcal{O}(\varepsilon^2)$
 $\phi_x(x, \pm \varepsilon h) = \phi_x(x, 0^\pm) \pm \varepsilon h \phi_{xy}(x, 0^\pm) + \mathcal{O}(\varepsilon^2)$ (12)

- Substitute asymptotic expansion:

$$\begin{aligned} \phi_y(x, \pm \varepsilon h) &\approx \varepsilon \phi_{1,y}(x, 0^\pm) \pm \varepsilon h \varepsilon \phi_{1,yy}(x, 0^\pm) \xrightarrow{\varepsilon^2 \ll 1} \\ \phi_x(x, \pm \varepsilon h) &\approx 1 + \varepsilon \phi_{1,x}(x, 0^\pm) \pm \varepsilon h \varepsilon \phi_{1,xy}(x, 0^\pm) \end{aligned} \quad (13)$$

- Substitute ⁽¹³⁾ into the BC (7):

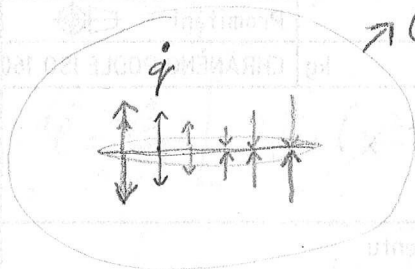
$$\varepsilon \phi_{1,y}(x, 0^\pm) = \pm \varepsilon h' \left[1 + \varepsilon \phi_{1,x}(x, 0^\pm) \right] \quad / : \varepsilon$$

$$\phi_{1,y}(x, 0^\pm) = \pm h' \quad (14)$$

discontinuous $\frac{\partial \phi_1}{\partial y} = 0$ at the airfoil

$\Rightarrow$ For the outer potential flow (excludes thin boundary layer) the airfoil is equivalent to a singular (infinitely thin) source/sink distribution at the axis.

$\nearrow \dot{Q} = 0$ zero net source $\Rightarrow$ mass is conserved



Flow described by (9, 10, 11, 14)