

# Subsonic thin profile - visualization

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```
[1]: import numpy as np  
import matplotlib.pyplot as plt
```

```
/home/lukas/.local/lib/python3.10/site-  
packages/matplotlib/projections/_init_.py:63: UserWarning: Unable to import  
Axes3D. This may be due to multiple versions of Matplotlib being installed (e.g.  
as a system package and as a pip package). As a result, the 3D projection is not  
available.  
    warnings.warn("Unable to import Axes3D. This may be due to multiple versions  
of "
```

## 1 Parameters

```
[2]: M      = 0.7          # Mach number  
epsilon = 0.1          # Profile thickness  
H      = lambda x: 0.5 - 2*x**2 # normalized thickness function  
dHdx   = lambda x: -4*x       # H'
```

## 2 Prandtl-Glauert transformation

```
[3]: beta   = np.sqrt(1-M**2)  
H_t    = lambda x: beta * H(x)  
dHdx_t = lambda x: beta * dHdx(x)
```

## 3 Analytical incompressible flow

The equivalent source/sink distribution  $q(x)$  along the axis of the profile is defined by the transformed asymptotic flow-tangency condition (6.3) as follows:

$$\begin{aligned} q(x) &= \bar{\phi}_{\bar{Y}}(X, 0^+) - \bar{\phi}_{\bar{Y}}(X, 0^-) \\ &= 2\bar{\phi}_{\bar{Y}}(X, 0^+) \\ &= 2\beta H'. \end{aligned}$$

The analytical solution for an incompressible potential flow driven by a continuous source sheet can be found, e.g., in [Schneider \(1978\)](#), p. 122:

$$\bar{\phi} = \frac{1}{2\pi} \int_{-1/2}^{1/2} q(\xi) \ln \sqrt{(X - \xi)^2 + \bar{Y}^2} d\xi.$$

Substituting the expression for  $q(x)$  we obtain

$$\bar{\phi} = \frac{\beta}{\pi} \int_{-1/2}^{1/2} H'(\xi) \ln \sqrt{(X - \xi)^2 + \bar{Y}^2} d\xi.$$

If we consider again the profile shape described by the quadratic function  $H(x) = 1/2 - 2X^2$ , then  $H' = -4X$  and

$$\bar{\phi} = \frac{-4\beta}{\pi} \int_{-1/2}^{1/2} \xi \ln \sqrt{(X - \xi)^2 + \bar{Y}^2} d\xi.$$

The analytical solution of the integral (found with Wolfram Mathematica) reads

$$\bar{\phi} = \frac{2\beta}{\pi} \left[ X - 2X\bar{Y} \left( \arctan \frac{1/2 - X}{\bar{Y}} + \arctan \frac{1/2 + X}{\bar{Y}} \right) + (1/4 - X^2 + \bar{Y}^2) \arctan \frac{4X}{1 + 4X^2 + 4\bar{Y}^2} \right].$$

More importantly, we are interested in the velocity components, that are defined as follows:

$$\begin{aligned} \bar{U} &= \frac{\partial \bar{\phi}}{\partial X} = \frac{1}{2\pi} \int_{-1/2}^{1/2} q(\xi) \frac{X - \xi}{(X - \xi)^2 + \bar{Y}^2} d\xi \\ &= \frac{\beta}{\pi} \int_{-1/2}^{1/2} H'(\xi) \frac{X - \xi}{(X - \xi)^2 + \bar{Y}^2} d\xi \\ &= \frac{-4\beta}{\pi} \int_{-1/2}^{1/2} \xi \frac{X - \xi}{(X - \xi)^2 + \bar{Y}^2} d\xi \end{aligned}$$

and

$$\begin{aligned}
\bar{V} &= \frac{\partial \bar{\phi}}{\partial \bar{Y}} = \frac{1}{2\pi} \int_{-1/2}^{1/2} q(\xi) \frac{\bar{Y}}{(X - \xi)^2 + \bar{Y}^2} d\xi \\
&= \frac{\beta}{\pi} \int_{-1/2}^{1/2} H'(\xi) \frac{\bar{Y}}{(X - \xi)^2 + \bar{Y}^2} d\xi \\
&= \frac{-4\beta}{\pi} \int_{-1/2}^{1/2} \xi \frac{\bar{Y}}{(X - \xi)^2 + \bar{Y}^2} d\xi.
\end{aligned}$$

Again, there exist analytical solutions of the integrals, and the velocity components can be expressed in closed form as follows:

$$\begin{aligned}
\bar{U} &= \frac{4\beta}{\pi} \left[ 1 - \bar{Y} \left( \arctan \frac{1/2 - X}{\bar{Y}} + \arctan \frac{1/2 + X}{\bar{Y}} \right) - X \operatorname{arctanh} \frac{4X}{1 + 4X^2 + 4\bar{Y}^2} \right] \\
\bar{V} &= \frac{-4\beta}{\pi} \left[ X \left( \operatorname{arccot} \frac{2\bar{Y}}{1 - 2X} + \arctan \frac{1/2 + X}{\bar{Y}} \right) + \frac{\bar{Y}}{2} \ln \left( 1 - \frac{8X}{(1 + 2X)^2 + 4\bar{Y}^2} \right) \right]
\end{aligned}$$

```
[4]: u_bar = lambda x,y_bar: 4*beta/np.pi * (
    -y_bar*( np.atan((0.5-x)/y_bar)
              +np.atan((0.5+x)/y_bar)
              )
    -x*np.atanh( 4*x/(
        1+4*x**2+4*y_bar**2
        )
    )
    )

v_bar = lambda x,y_bar: -4*beta/np.pi *( x*( np.atan((1-2*x)/(2*y_bar))
                                             +np.atan((1+2*x)/(2*y_bar))
                                             )
    +y_bar/2 *np.log( 1 - 8*x/(
        (1+2*x)**2
        +4*y_bar**2
        )
    )
    )
```

Thus, we can plot the auxiliary incompressible flow over the transformed profile defined by  $\bar{Y} = \pm\epsilon\beta H(X)$ .

### 3.1 Visualization

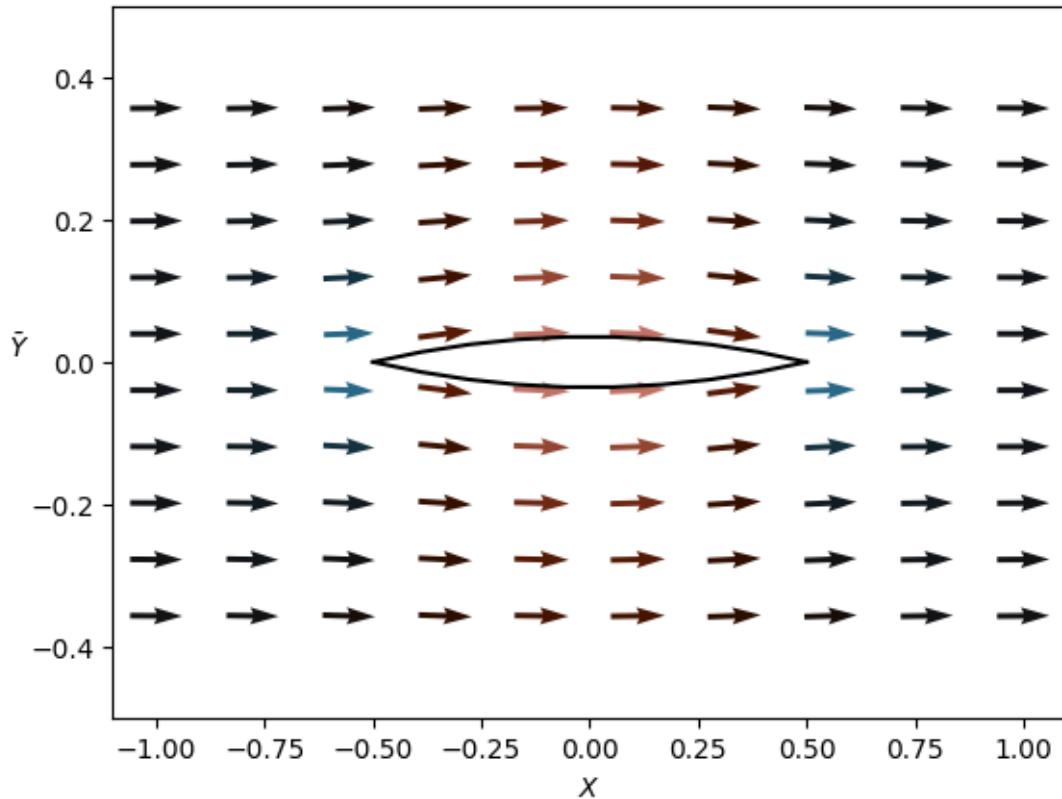
```
[5]: xmin, xmax = -1.0 , 1.0 # boundaries of the visualized region  
ymin, ymax = -0.5 , 0.5  
N           = 10          # number of velocity vectors in each direction
```

```
[6]: x,y = np.meshgrid( np.linspace(xmin,xmax,N) ,  
                      np.linspace(ymin,ymax,N) )  
y_bar = beta * y
```

```
[7]: u_ = 1 + epsilon * u_bar(x,y_bar)  
v_ =      epsilon * v_bar(x,y_bar)
```

```
[8]: # Velocity vectors  
plt.quiver(x,y_bar,u_,v_,np.sqrt(u_**2+v_**2),angles='xy',  
            pivot='mid',clim=(1-epsilon,1+epsilon),cmap='berlin')  
  
# Shape of the profile  
x_shape      = np.linspace(-0.5, 0.5, N)  
y_shape_bar = epsilon * beta * H(x_shape)  
plt.plot(x_shape, y_shape_bar, 'k-', x_shape, -y_shape_bar, 'k-')  
  
# Labels  
plt.xlabel(r'$X$')  
plt.ylabel(r'$\bar{Y}$',rotation='horizontal')  
plt.ylim(ymin, ymax)
```

```
[8]: (-0.5, 0.5)
```



## 4 Compressible flow

The subsonic compressible flow is obtained by backward transformation into the original variables

$$Y = \bar{Y}/\beta$$

$$U_1 = \bar{U}/\beta^2$$

$$V_1 = \bar{V}/\beta$$

```
[9]: u1 = lambda x,y: u_bar(x,y*beta) / beta**2
v1 = lambda x,y: v_bar(x,y*beta) / beta
```

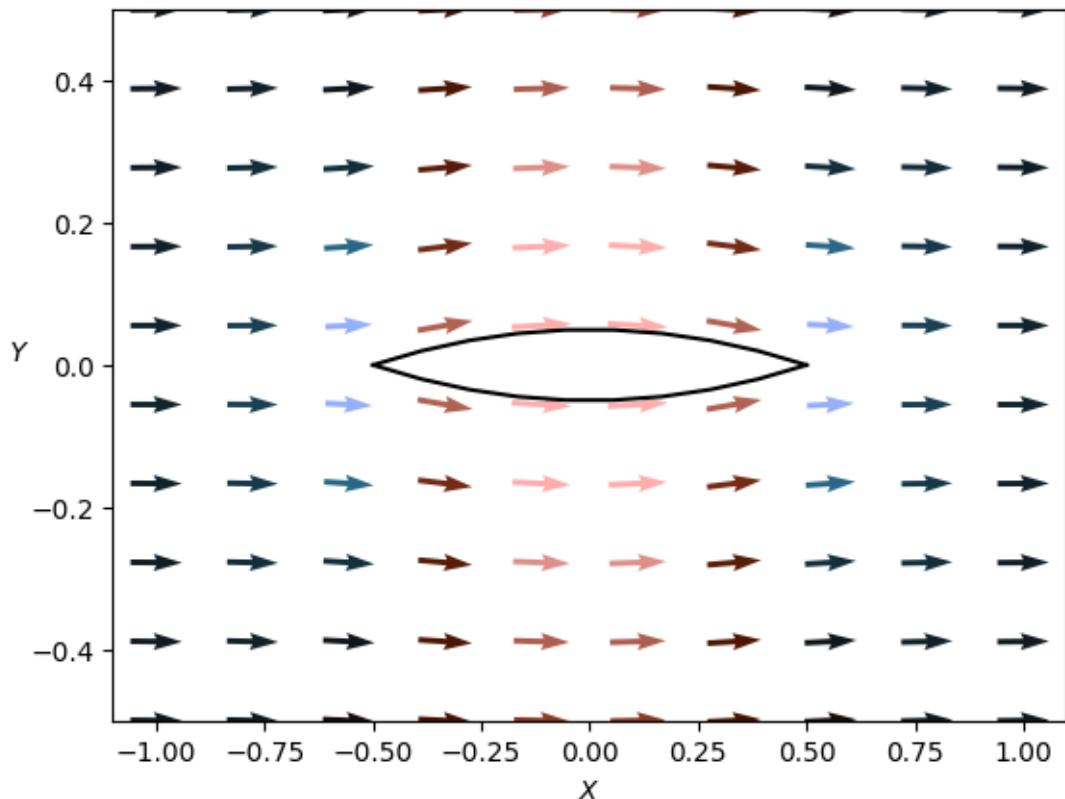
```
[10]: u = 1 + epsilon * u1(x,y)
v = epsilon * v1(x,y)

y_shape = y_shape_bar / beta
```

```
[11]: plt.quiver(x,y,u,v,np.sqrt(u**2+v**2),angles='xy',
pivot='mid',clim=(1-epsilon,1+epsilon),cmap='berlin')
plt.plot(x_shape, y_shape, 'k-', x_shape, -y_shape, 'k-')
```

```
# Labels  
plt.xlabel(r'$X$')  
plt.ylabel(r'$Y$', rotation='horizontal')  
plt.ylim(ymin, ymax)
```

[11]: (-0.5, 0.5)



[ ]: